

Analysis of Pressure-Driven Heat and Mass Transfer of Hydromagnetic Flow Past Darcy-Forchheimer Porous Media Using Lie Group

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Abstract: This study investigates heat and mass transfer of Magnetohydrodynamic (MHD) pressure-driven flow in a Darcy-Forchheimer porous medium with inclined uniform magnetic field and thermal radiation. The governing equations of the model are reduced to a system of coupled non-linear ordinary differential equations by applying a scaling group of transformations. The coupled differential equations are then solved using Weighted Residual Method (WRM) with the results compared with that obtained from shooting technique coupled with fourth order Runge-Kutta method. The obtained results are presented graphically to represent the effect various fluid parameters on the dimensionless velocity, temperature, concentration and pressure drop. Finally, the effects of Skin friction, Nusselt and Sherwood numbers which are of physical and engineering interest are presented and discussed.

Key words: Magnetohydrodynamic, pressure-driven flow, porous medium, Darcy-Forchheimer, weighted residual method, thermal radiation

INTRODUCTION

The study of MHD fluid flow induced by the concurrent actions of buoyancy forces follow-on from the mass diffusion and thermal conductivity is gaining attention due to its interesting applications. The radiative thermal physics of MHD problems with heat and mass transfer through Darcy-Forchheimer permeable medium are usually come across in the power engineering and other pollutant, solidification and evaporation cooling, grain storage are areas in which the combined thermo-solutal in a permeable medium are experienced.

In few of its applications, Singh *et al.* (2012) examined heat and mass transfer in MHD flow through a vertical stretching surface in the occurrence of heat generation. Analysis of the MHD stagnation point fluid flow with the effects of chemical reaction on heat and mass transfer near a permeable stretching vertical sheet and injection or suction by Lie group was investigated by Afify and Elgazery (2012). The result presents monotonically decreased when the magnetic field parameter was increased. Alam *et al.* (2003) reported on viscous dissipative effect of MHD convective flow over an inclined permeable stretching sheet and chemical reaction. It was found that chemical reaction may significantly influenced the concentration boundary layer and wall mass transfer rate while Ahmed and Bhattacharyya (2014) studied free convection chemical reaction in MHD couette

flow with heat source. The above researchers did not consider the effects of pressure gradient, porosity and radiation on the fluid flow.

The analysis of flow past porous medium has received the attention of several researchers due to its extensive applications in enhancing recovery of petroleum, chemical engineering, etc. The thermal diffusion and heat generation of heat and mass transfer in MHD fluid flow through a vertical motioning permeable sheet and medium was examined by Saxena and Dubey (2011) while Hunegnaw and Kishan (2014) examined heat and mass transfer in a MHD flow along a moving surface with chemical reaction, variable properties in a porous medium and viscous dissipation. The cited researchers did not consider the significant of non-Darcy porous medium and radiation on the flow as well as the influence of fluid parameters on the pressure.

In studies of heat and mass transfer, Ishak and Nazar (2006) analyzed boundary layers of a permeable medium over a vertical surface by applying Brinkman equation for the porous medium model while Loganathan and Arasu (2010) investigated thermophoresis influences on heat and mass transfer in a MHD Darcy-Forchheimer permeable medium in the occurrence injection/suction over a porous wedge. It was noticed that the flow velocity decreased as the inertial term increased, the behavior was due to the fact that the inertia parameter supports the resistance of the flow fluid mechanism. Also,

Ahammad *et al.* (2013) examined heat and mass transfer through a non-Darcy permeable medium in the presence of magnetic field and (Hemalatha *et al.*, 2014) investigated thermal dispersion with melting on Darcy-Forchheimer porous medium in a hydromagnetic flow over a stretching surface. Forchheimer extension for the flow equations in steady state was considered.

Moreover, Lee *et al.* (2008) examined radiative heat transfer past a porous sheet using Lie group. It was obtained that the heat and velocity boundary layer thicknesses increased by a rise in the radiation term, variational increased in the Prandtl number slow down both the temperature and velocity in the boundary layer. Velocity increased and temperature decreased when there was an increased in Grashof number and the porosity parameter. Chauhan and Rastogi (2010) analyzed natural convection effects on radiative MHD flow through a vertical partially filled porous medium without giving attention to the effect of mass diffusion and heat generation in the studied. Dada and Adefolaju (2012) reported on hydrodynamic flow fluid over a boundary layer in a stretching sheet entrenched in a radiation and non-Darcy permeable medium. It was found that an increase in the Forchheimer number reduced the velocity profile. However, the study did not includes the effect of pressure gradient, heat source and chemical reaction rate on the flow while Kumar and Sivaraj (2012) carried out analysis on viscoelastic fluid of MHD flow in non-Darcy permeable medium along a motioning vertical conduit.

Follow from the above studies, the researchers neglected the influence of Darcy-Forchheimer permeable medium inclined magnetic field and thermal radiation on a fluid flow as well as the effects of fluid parameters on the pressure. The current study examine the combined influences of inclined magnetic field, pressure drop, thermal radiation and Darcy-Forchheimer porous medium in a steady convective heat and mass transfer in a MHD flow. The fluid flow is pressure-driven past a permeable plate with inclined uniform magnetic field.

MATERIALS AND METHODS

Formulation of the problem: Examine heat and mass transfer of MHD pressure-driven, steady, viscous, laminar and incompressible fluid over a permeable stretching plate in Darcy-Forchheimer medium and radiation under the effect of inclined magnetic field and pressure gradient. The fluid motion is maintained by both pressure gradient and gravity and the flow is to be in the direction of X and Y-axis normal to it. A uniform magnetic field B_0 is introduced at angle α lying in the range $0 < \alpha < \pi/2$ in the

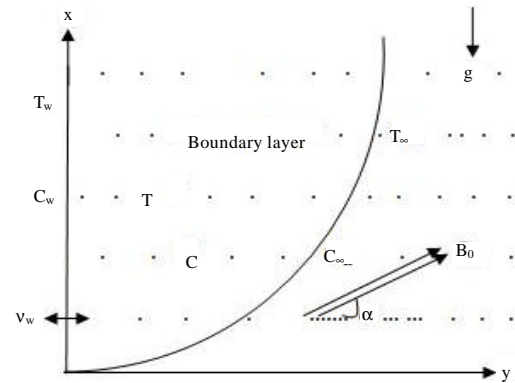


Fig. 1: Geometry of the problem

direction of the flow. The plate is maintained at the temperature and species concentration T_w , C_w and free stream temperature and species concentration T_∞ , C_∞ , respectively. The equations governing the steady radiative heat and mass transfer of two-dimensional magnetohydrodynamics pressure-driven fluid flow in Darcy Forchheimer porous medium with inclined magnetic field are as follows (Fig. 1):

$$U \frac{\partial U}{\partial X} + V \frac{\partial U}{\partial Y} = 0 \quad (1)$$

$$U \frac{\partial U}{\partial X} + V \frac{\partial U}{\partial Y} = -\frac{1}{\rho} \sigma B_0^2 U \sin^2 \alpha - \frac{1}{\rho} \frac{\partial P}{\partial X} + \nu \left(\frac{\partial^2 U}{\partial X^2} + \frac{\partial^2 U}{\partial Y^2} \right) - \frac{\nu}{K^*} U - \frac{b}{K^*} U^2 + g \beta_T (T - T_\infty) + g \beta_C (C - C_\infty) \quad (2)$$

$$U \frac{\partial V}{\partial X} + V \frac{\partial V}{\partial Y} = -\frac{1}{\rho} \frac{\partial P}{\partial Y} + \nu \left(\frac{\partial^2 V}{\partial X^2} + \frac{\partial^2 V}{\partial Y^2} \right) \quad (3)$$

$$\rho C_p \left(U \frac{\partial T}{\partial X} + V \frac{\partial T}{\partial Y} \right) = k \left(\frac{\partial^2 T}{\partial X^2} + \frac{\partial^2 T}{\partial Y^2} \right) - \frac{\partial q_x}{\partial X} - \frac{\partial q_y}{\partial Y} + Q_0 (T - T_\infty) \quad (4)$$

$$U \frac{\partial C}{\partial X} + V \frac{\partial C}{\partial Y} = D \left(\frac{\partial^2 C}{\partial X^2} + \frac{\partial^2 C}{\partial Y^2} \right) - \gamma (C - C_\infty) \quad (5)$$

The corresponding initial and boundary conditions are as follows:

$$\begin{aligned} U = 0, V = v_w, P = 0, T = T_\infty + (T_w - T_\infty)AX, \\ C = C_\infty + (C_w - C_\infty)BX \text{ at } Y = 0 \\ U = 0, T = T_\infty, C = C_\infty \text{ as } Y \rightarrow \infty \end{aligned} \quad (6)$$

where, U , V , P , C and T are the velocity component in the X direction, velocity component in the Y direction, pressure, concentration of species in the fluid and temperature of the fluid, respectively. A and B are constants defined as $A = B = 1/l$, B_0 is the magnetic field strength, α is the angle of inclination of the magnet, $v_w > 0$ indicates wall injection and $v_w < 0$ indicates wall suction, respectively. The physical quantities ν , b , K^* , ρ , α , D , C_p , k , Q_0 and γ are the fluid kinematics viscosity, Forchheimer parameter of the medium, permeability of the porous medium, density, electric conductivity of the fluid, mass diffusion coefficient, specific heat at constant pressure, thermal conductivity, internal heat generation or absorption and chemical reaction coefficient, respectively. g is the gravitational acceleration, β_T and β_C are the thermal and concentration expansion coefficients, respectively. q_x and q_y are the radiative heat flux in the X and Y direction, respectively. Using Rosseland diffusion approximation for radiation by Kareem *et al.* (2018):

$$q_x = -\frac{4\sigma_0}{3\delta} \frac{\partial T^4}{\partial X} \text{ and } q_y = -\frac{4\sigma_0}{3\delta} \frac{\partial T^4}{\partial Y} \quad (7)$$

where, σ_0 and δ are the Stefan-Boltzmann and the mean absorption coefficient, respectively. Assume the temperature difference within the flow are sufficiently small such that T^4 may be expressed as a linear function of temperature, using Taylor series to expand T^4 about the free stream T_∞ and neglecting higher order terms, this gives the approximation:

$$T^4 \approx 4T_\infty^3 T - 3T_\infty^4 \quad (8)$$

Using Eq. 8 then, Eq. 7 can be express as:

$$\frac{\partial q_x}{\partial X} = -\frac{16\sigma_0 T_\infty^3}{3\delta} \frac{\partial^2 T}{\partial X^2} \text{ and } \frac{\partial q_y}{\partial Y} = -\frac{16\sigma_0 T_\infty^3}{3\delta} \frac{\partial^2 T}{\partial Y^2} \quad (9)$$

Introducing the following non-dimensional quantities:

$$\begin{aligned} x = \frac{X}{l}, y = \frac{Y}{l}, u = \frac{Ul}{\nu}, v = \frac{Vl}{\nu}, p = \frac{Pl^2}{\rho\nu^2}, \\ \phi = \frac{C - C_\infty}{C_w - C_\infty}, \theta = \frac{T - T_\infty}{T_w - T_\infty} \end{aligned} \quad (10)$$

Substituting (Eq. 9) and (Eq. 10) into Eq. 1-6, to obtain:

$$u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial y} = 0 \quad (11)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -H_a^2 \sin^2 \alpha u - \frac{\partial p}{\partial x} + \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) - D_a u - F_s u^2 + G_r \theta + G_c \phi \quad (12)$$

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{\partial p}{\partial y} + \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) \quad (13)$$

$$u \frac{\partial \theta}{\partial x} + v \frac{\partial \theta}{\partial y} = \frac{1}{Pr} \left(1 + \frac{4}{3} R \right) \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} \right) + Q \theta \quad (14)$$

$$u \frac{\partial \phi}{\partial x} + v \frac{\partial \phi}{\partial y} = \frac{1}{Sc} \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) - \lambda \phi \quad (15)$$

The corresponding initial and boundary conditions are follows:

$$\begin{aligned} u = 0, v = -f_w, p = 0, \theta = x, \phi = x \text{ at } y = 0 \\ u = 0, \theta = 0, \phi = 0 \text{ as } y \rightarrow \infty \end{aligned} \quad (16)$$

where, the term $Pr = \mu C_p / k$ is the Prandtl, $Sc = \nu / D$ is the Schmidt, $H_a = B_0 \sqrt{\sigma / \mu}$ is the Hartmann, $G_r = l^3 g \beta_T (T_w - T_\infty) / \nu^2$ is the thermal Grashof, $G_c = l^3 g \beta_C (C_w - C_\infty) / \nu^2$ is the solutal Grashof, $Q = l^2 Q_0 / \mu C_p$ is the heat source, $\lambda = l^2 \gamma / \nu$ is the concentration parameter. $D_a = l^2 / K^*$ is the Darcy parameter, $F_s = la / K^*$ is the Forchheimer inertia term, $R = 4\sigma_0 T_\infty^3 / \delta k$ is the Radiation heat transfer parameter and $f_w = v_w l / \nu$ is the non-dimensional wall mass transfer coefficient, respectively.

Introducing the stream function, $u = \partial \psi / \partial y$, $v = \partial \psi / \partial x$ continuity equation is automatically satisfied and Eq. 12-16 becomes:

$$\frac{\partial \psi}{\partial y} \frac{\partial^2 \psi}{\partial x \partial y} - \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial y^2} = -H_a^2 \sin^2 \alpha \left(\frac{\partial \psi}{\partial y} \right) - \frac{\partial p}{\partial x} + \left(\frac{\partial^3 \psi}{\partial x^2 \partial y} + \frac{\partial^3 \psi}{\partial y^3} \right) - D_a \left(\frac{\partial \psi}{\partial y} \right) - F_s \left(\frac{\partial \psi}{\partial y} \right)^2 + G_r \theta + G_c \phi \quad (17)$$

$$-\frac{\partial \psi}{\partial y} \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial x \partial y} = -\frac{\partial p}{\partial y} - \left(\frac{\partial^3 \psi}{\partial x^3} + \frac{\partial^3 \psi}{\partial x \partial y^2} \right) \quad (18)$$

$$\frac{\partial \psi}{\partial y} \frac{\partial \theta}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial \theta}{\partial y} = \frac{1}{P_r} \left(1 + \frac{4}{3} R \right) \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} \right) + Q \theta \quad (19)$$

$$\frac{\partial \psi}{\partial y} \frac{\partial \phi}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial \phi}{\partial y} = \frac{1}{S_c} \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) - \lambda \phi \quad (20)$$

Subject to the initial and boundary conditions:

$$\begin{aligned} \frac{\partial \psi}{\partial y} = 0, \frac{\partial \psi}{\partial x} = f_w, p = 0, \theta = x, \phi = x \text{ at } y = 0 \\ \frac{\partial \psi}{\partial y} = 0, \theta = 0, \phi = 0 \text{ as } y \rightarrow \infty \end{aligned} \quad (21)$$

Introducing simplified form of Lie-group transformations namely, the scaling group of transformations to Eq. 17-21 which is equivalent to determining the invariant solutions of these equations under a continuous one-parameter group according to Pramanik (2013). One of the methods is to search for a transformation group from an elementary set of one-parameter scaling group of transformations, given as ∇ :

$$\begin{aligned} \nabla: x^* = x e^{\epsilon \alpha_1}, y^* = y e^{\epsilon \alpha_2}, \psi^* = \psi e^{\epsilon \alpha_3}, \\ u^* = u e^{\epsilon \alpha_4}, v^* = v e^{\epsilon \alpha_5}, \\ p^* = p e^{\epsilon \alpha_6}, \theta^* = \theta e^{\epsilon \alpha_7}, \phi^* = \phi e^{\epsilon \alpha_8} \end{aligned} \quad (22)$$

where, α_1 - α_8 are transformation parameters of the group to be determined later and ϵ is a small parameters. The task is to find relationships among the exponents α 's such that (Eq. 17-21) will remain invariant under the point transformations. Substituting transformation (Eq. 22) into Eq. 17-21 and applying invariant conditions yields the similarity transformations:

$$\begin{aligned} \eta = y, \psi = x f(\eta), p = p_d(\eta), \\ \theta = x \theta(\eta), \phi = x \phi(\eta) \end{aligned} \quad (23)$$

Substituting, the similarity variables (Eq. 23) into Eq. 17-21. The following system of non-linear differential equations are obtained:

$$f'' + f f' - (1 + F_s) f'^2 - (H_a^2 \sin^2 \alpha + D_a) \quad (24)$$

$$\begin{aligned} f' + G_r \theta + G_c \phi = 0 \\ -p_d' = f' + f f' \end{aligned} \quad (25)$$

$$\left(1 + \frac{4}{3} R \right) \theta'' + P_r \theta' - P_r f' \theta + P_r Q \theta = 0 \quad (26)$$

$$\phi'' + S_c f \phi' - S_c f' \phi - S_c \lambda \phi = 0 \quad (27)$$

The corresponding initial and boundary conditions take the form:

$$\begin{aligned} f = f_w, f' = 0, p_d = 0, \theta = 1, \phi = 1 \text{ at } \eta = 0 \\ f' = 0, \theta = 0, \phi = 0 \text{ as } \eta \rightarrow \infty \end{aligned} \quad (28)$$

Integrating, Eq. 27 with the initial and boundary conditions when $f_w = 1$ let pressure drop $-p_d = G$, this become:

$$G = f' + \frac{1}{2} f^2 - \frac{1}{2} \quad (29)$$

Applying WRM to Eq. 26-29 by assuming a polynomial of with unknown coefficients or parameters to be determined later, this polynomial is called the trial function:

$$f(\eta) = \sum_{i=0}^n a_i \eta^i, \theta(\eta) = \sum_{i=0}^n b_i \eta^i, \phi(\eta) = \sum_{i=0}^n c_i \eta^i \quad (30)$$

Impose the boundary conditions (Eq. 28) on the trial functions also substituting the trial functions into Eq. 24, 26 and 27 to obtain the residual equations. Minimizing the residual error by forcing it to zero at some set of collocation points within the domain in order to obtain the unknown coefficients using Maple 2016 software:

$$\begin{aligned} f(\eta) = 1.000000000 + 1.352521223 \eta^2 - \\ 2.136719818 \eta^3 + 2.081355570 \eta^4 - \\ 1.497172413 \eta^5 + 0.8258367187 \eta^6 - \\ 0.3501466456 \eta^7 + 0.1122038542 \eta^8 - \\ 0.02624357625 \eta^9 + 0.004215585400 \eta^{10} - \\ 0.0004145534793 \eta^{11} + \\ 0.00001876803320 \eta^{12} \end{aligned} \quad (31)$$

$$\begin{aligned} \theta(\eta) = 1.000000000 - 0.5182224615 \eta + \\ 0.003827280553 \eta^2 + 0.1483752625 \eta^3 - \\ 0.1835687527 \eta^4 + 0.1537432440 \eta^5 - \\ 0.09560654092 \eta^6 + 0.04481892093 \eta^7 - \end{aligned}$$

$$0.01564765871\eta^8 + 0.003938812348\eta^9 - \\ 0.0006738700894\eta^{10} + 0.00006996577184\eta^{11} - \\ 0.000003320194742\eta^{12}$$

$$\phi(\eta) = 1.000000000 - 1.181989422\eta + \\ 0.5262268881\eta^2 + 0.09051904287\eta^3 - \\ 0.3607131506\eta^4 + 0.3626303004\eta^5 - \\ 0.2451914515\eta^6 + 0.1220539007\eta^7 - \\ 0.04472607833\eta^8 + 0.01171007571\eta^9 - \\ 0.002067978267\eta^{10} + 0.0002202548424\eta^{11} - \\ 0.00001066883171\eta^{12}$$

Differentiate Eq. 31 to obtain:

$$f'(\eta) = 2.705042446\eta - 6.410159454\eta^2 + \\ 8.325422280\eta^3 - 7.485862065\eta^4 + \\ 4.955020312\eta^5 - 2.451026519\eta^6 + \\ 0.8976308336\eta^7 - 0.2361921862\eta^8 + \\ 0.04215585400\eta^9 - 0.004560088272\eta^{10} + \\ 0.0002252163984\eta^{11}$$

Also, substituting for f and f' in Eq. 29 with the corresponding constant values to obtain the pressure drop as:

$$G(\eta) = -0.5000000000 + 2.705042446\eta - \\ 6.410159454\eta^2 + 8.325422280\eta^3 - \\ 7.485862065\eta^4 + 4.955020312\eta^5 - \\ 2.451026519\eta^6 + 0.8976308336\eta^7 - \\ 0.2361921862\eta^8 + 0.04215585400\eta^9 - \\ 0.004560088272\eta^{10} + 0.0002252163984\eta^{11} + \\ \frac{1}{2}(1.000000000 + 1.357521223\eta^2 - \\ 2.136719818\eta^3 + 2.081355570\eta^4 - \\ 1.497172413\eta^5 + 0.8258367187\eta^6 - \\ 0.3501466456\eta^7 + 0.1122038542\eta^8 - \\ 0.02624357625\eta^9 + 0.004215585400\eta^{10} - \\ 0.0004145534793\eta^{11} + 0.00001876803320\eta^{12})^2$$

The physical quantity of practical interest are the local skin friction C_f , the Nusselt number Nu and the local Sherwood number Sh define as:

$$C_f = \frac{\tau_w}{\rho u_w^2}, \quad Nu = \frac{q_w x}{k(T_w - T_\infty)}, \quad Sh = \frac{q_m x}{D(C_w - C_\infty)} \quad (36)$$

where, k is the thermal conductivity of the fluid, τ_w , q_w and q_m are respectively given by:

$$\tau_w = -\mu \left(\frac{\partial u}{\partial y} \right)_{y=0}, \quad q_w = -k \left(\frac{\partial T}{\partial y} \right)_{y=0}, \quad q_m = -D \left(\frac{\partial C}{\partial y} \right)_{y=0} \quad (37)$$

Therefore, the local skin friction coefficient, local Nusselt number and local Sherwood number are:

$$C_f Re_x^{\frac{1}{2}} = -f'(0), \quad Nu Re_x^{\frac{1}{2}} = -\theta'(0), \quad Sh Re_x^{\frac{1}{2}} = -\phi'(0) \quad (38)$$

where, $Re_x = v_\infty x / \nu$ is the Reynolds number. The process of weighted residual method are repeated for different values of G_r , G_c , H_a , α , Q , P_r , D_a , R , F_s , S_c and λ . The following computational results in the tables were obtained and compared with Runge-Kutta method generated by maple software.

RESULTS AND DISCUSSION

In order to get clear insight of the physical problem, numerical computations has been carried out using the method of Weighted Residual for the representative velocity field, pressure field, temperature field, concentration field, coefficient of the skin friction, the rates of heat and mass transfer in terms of Nusselt number and Sherwood number, respectively at the plate have been carried out for different values of the parameters. The following default parameter values are adopted for computation: $G_r = G_c = 5$, $Q = \lambda = 0.5$, $R = F_s = D_a = 1$, $P_r = 0.72$, $S_c = 0.62$ and $M = 3$ at $H_a = 6$ and $\alpha = 30^\circ$. All graphs, therefore, correspond to these values unless specifically indicated on the appropriate graph.

Figure 2 and 3 bring out clearly the effect of the Hartmann number H_a on the velocity and pressure

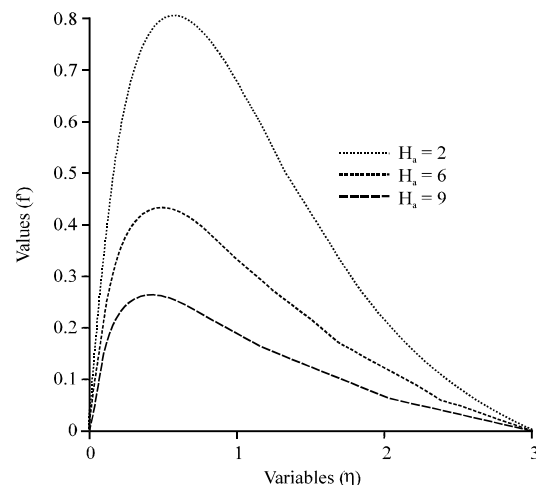


Fig. 2: Velocity profiles for different values of H_a

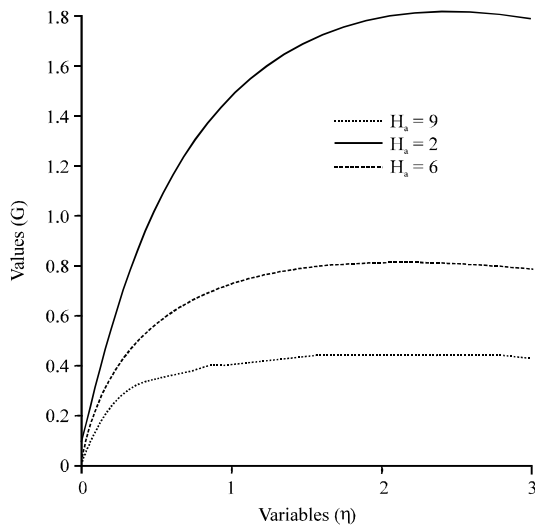


Fig. 3: Pressure profiles for different values of H_a

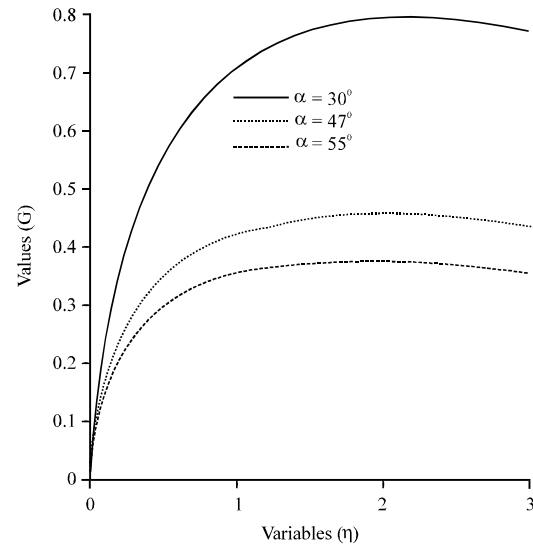


Fig. 5: Pressure profiles for different values α

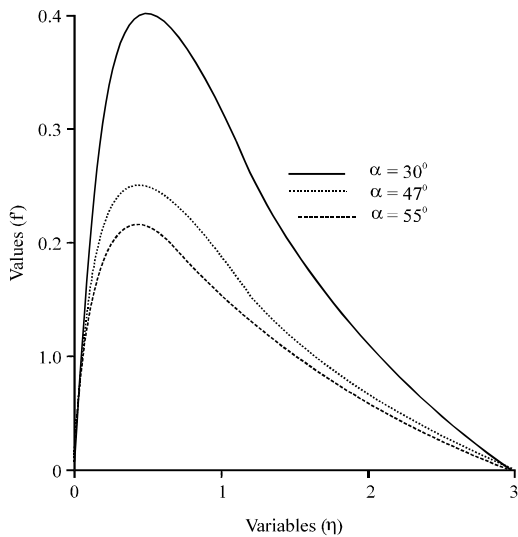


Fig. 4: Velocity profiles for different values α

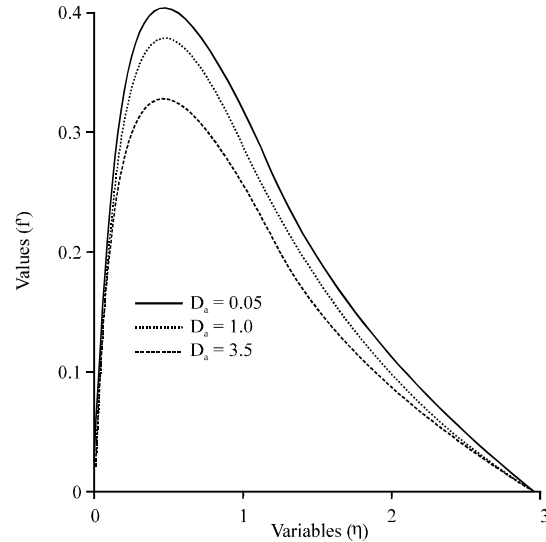


Fig. 6: Velocity profiles for different values D_a

boundary layer thickness, since, the magnetic field exerts a retarding force on the free convection fluid flow. Increasing H_a decreases the velocity and pressure boundary layer.

Figure 4 and 5 show the velocity and pressure profiles for different angles of inclination of the magnetic field α , an increase in the angle of inclination increases the effect of the buoyancy force and consequently, the driving force to the fluid flow decreases as a result, velocity and pressure boundary layer thickness decreases.

Figure 6-9 show the effect of the porosity parameter D_a and inertial parameter F_s on the velocity and pressure profiles. It was observed that the velocity and pressure

decreases as the porosity and inertial parameter increases. The reason for this behavior is that the wall of the surface provides an additional resistance to the fluid flow mechanism which causes the fluid to move at a retarded rate.

Figure 10 and 11 depict the influence of different values of the Prandtl number P_r on the pressure and temperature distribution. It is observed that an increase in the ratio of momentum diffusivity to thermal diffusivity results in the, respectively decrease in pressure and temperature profiles. The reason for this behavior is that an increase in the P_r results in a decrease in the boundary layer thickness and reduce the average temperature within the boundary layer because smaller values of P_r are

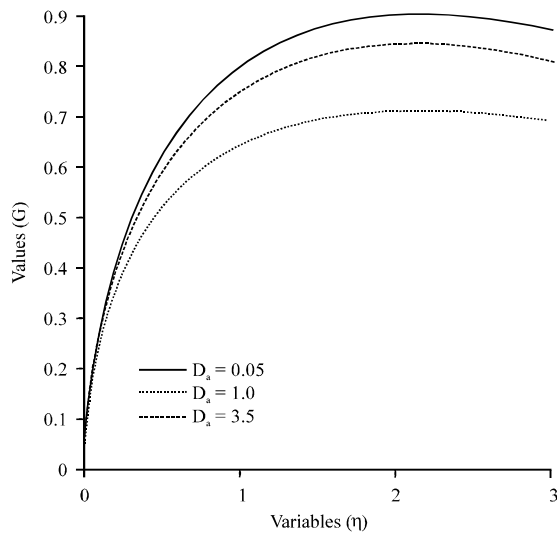


Fig. 7: Pressure profiles for different values of D_a

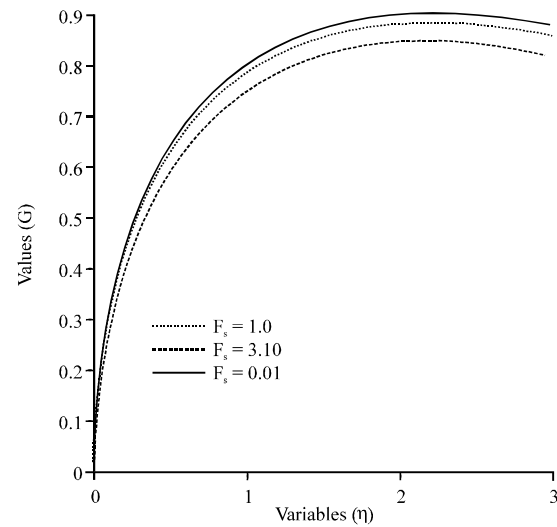


Fig. 9: Pressure profiles for different values of F_s

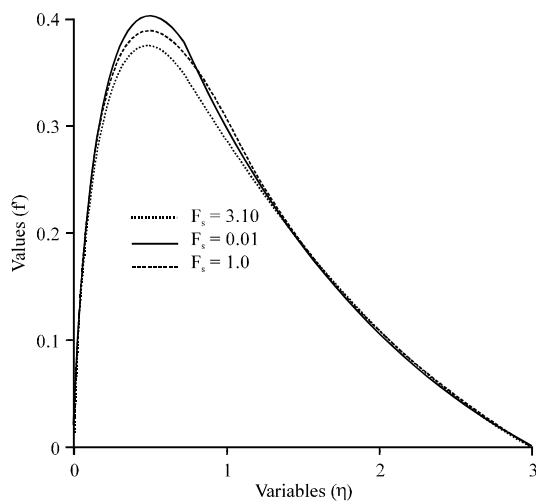


Fig. 8: Velocity profiles for different values of F_s

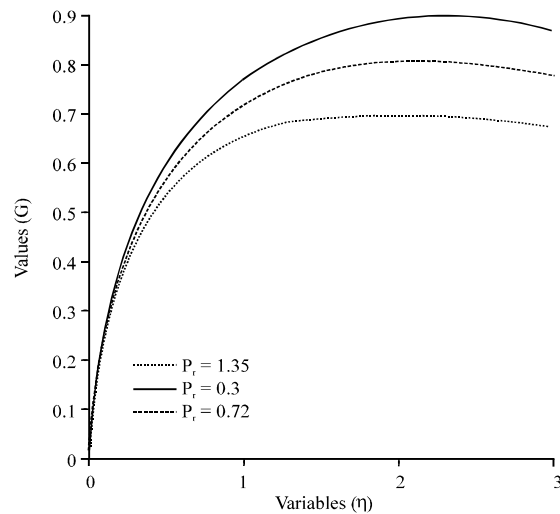


Fig. 10: Pressure profiles for different values of P_r

equivalent to increasing in the thermal conductivities, therefore, heat is able to diffuse away from the heated plate than higher values of P_r .

The effect of Schmidt number S_c on the pressure and concentration profiles are represented on Fig. 12 and 13. Schmidt number is the ratio of the momentum to the mass diffusivity. An increase in S_c causes reductions in the pressure and concentration profiles which are accompanied by simultaneous decrease in the pressure and concentration boundary layers. Schmidt number is therefore, quantifies the relative effectiveness of momentum and mass transport by diffusion in the hydrodynamic pressure and concentration boundary layers. All these agrees with the expectations.

Table 1 depicts the effect of some physical parameters on skin friction, nusselt and Sherwood number. It clearly shows that, the magnetic term H_s and angle of inclination α have a retarding effect on the skin friction, nusselt and Sherwood number, respectively. Increase in heat source increase skin friction and Sherwood number while it decreases the Nusselt number because heat within the boundary layer reduces. Also, number Prandtl number P_r has decreasing effect on the skin friction and Sherwood number but has an increasing effect on Nusselt number because of the thickness of the thermal boundary layer. Moreover, skin friction and nusselt have a decelerating effect on increasing Schmidt number but an increasing effect on increasing Schmidt number.

Table 1: Effect of H_a , α , Q , P_r , S_c on τ , Nu and Sh (PP-Physical Parameters)

PP	Values	Weighted Residual method			4th order R-K		
		τ	Nu	Sh	τ	Nu	Sh
H_a	2	4.05305	0.64028	1.31465	4.05634	0.64017	1.31431
	6	2.70504	0.51822	1.18199	2.71947	0.51806	1.18134
	9	2.00826	0.45420	1.10920	1.98953	0.45406	1.10844
α	30°	2.70504	0.51822	1.18199	2.71947	0.51807	1.18134
	47°	2.05511	0.45787	1.11345	2.02117	0.45772	1.11269
	55°	1.86538	0.44189	1.09484	1.82722	0.44175	1.09408
Q	0.5	2.70504	0.51822	1.18199	2.71947	0.51806	1.18134
	0.8	2.74628	0.42937	1.18688	2.74405	0.42923	1.18622
	1.2	2.78445	0.29558	1.19442	2.79831	0.29547	1.19377
P_r	0.3	2.75338	0.40518	1.19053	2.76763	0.40513	1.18987
	0.72	2.70504	0.51822	1.18199	2.71947	0.51806	1.18134
	1.35	2.63053	0.71094	1.16950	2.64527	0.71056	1.16886
S_c	0.4	2.79921	0.52961	0.90598	2.79305	0.52945	0.90565
	0.62	2.705042	0.518222	1.181989	2.71947	0.51806	1.18134
	1.5	2.472478	0.493724	2.173529	2.48940	0.49357	2.17044

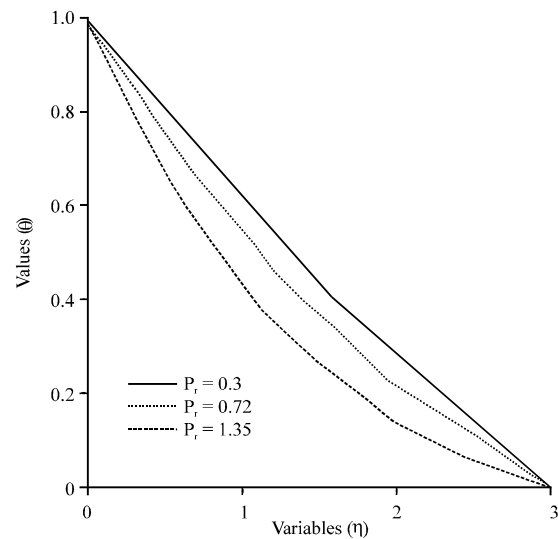


Fig. 11: Temperature profiles for different values of P_r

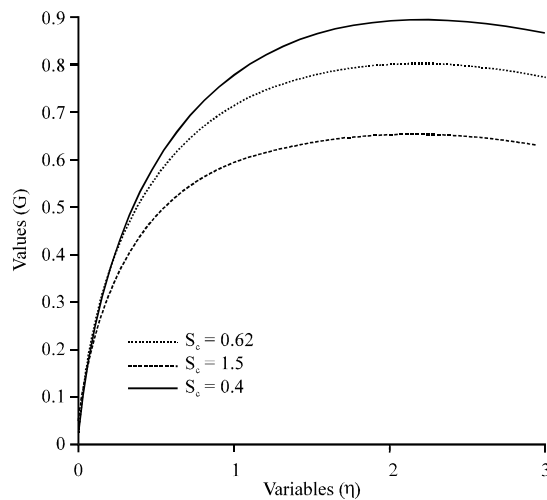


Fig. 12: Pressure profiles for different values of S_c

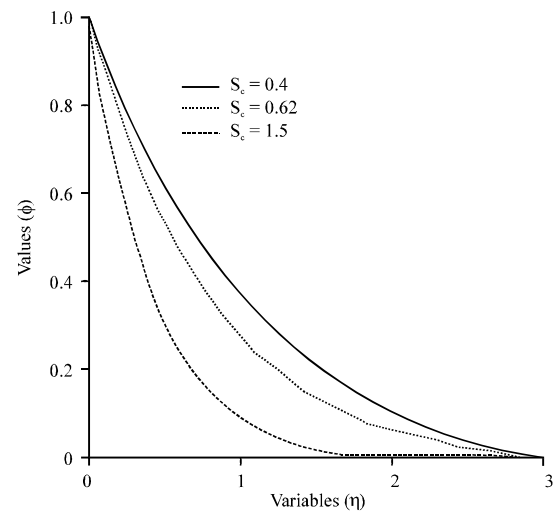


Fig. 13: Concentration profiles for different values of S_c

CONCLUSION

The non-dimensional of the formulated governing equations are reduced to a couple ordinary differential equations by using scaling group transformation. The numerical solution are obtained using Weighted Residual method. From the numerical results, it is seen that, an increase in the values of Hartmann, degree of inclination of the magnetic field, porosity parameter, inertial parameter, Prandtl or Schmidt numbers is manifested as a decrease in the flow velocity, pressure, temperature and concentration distribution.

The results of the study are interested because flows past a permeable surface play an essential role in applications of science and engineering, also in many transport processes in nature.

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