

A New Method for Cylindrical Multilayered Waveguides Analysis

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Abstract: Cylindrical multilayered waveguides are often used as parts of various microwave and optical devices. Layered structure of waveguide filling allows more flexibility of tuning and optimization of the device parameters. In the waveguide with inhomogeneous cross-section, one can obtain order of cut-off frequencies of eigen modes which is different from the normal one. The study describes a new technique of obtaining the dispersion characteristics of such waveguides based on tensor Green's functions method. The proposed method is effective in constructing the dispersion equations as it uses microwave circuits theory description and recurrent expressions instead of direct solution of the system of boundary equations.

Key words: Waveguides, propagation constant, transmission lines, Green's function methods, analysis

INTRODUCTION

Waveguides are the basic structures for microwave resonators, filters, rotary joints, etc. Waveguides with inhomogeneous filling can be used to obtain, e.g., low dispersion and attenuation in transmission line, high-Q resonators or can serve as parts of mode filters. The particular form of these filling is radially piecewise inhomogeneous cylindrical waveguides such as circular waveguides, coaxial lines, optical fibers and Goubau lines. One of the main characteristics of such structures is a propagation constant (wave number) which contains the information about phase constant and attenuation. The most popular for analysis of multilayered waveguides is the method in which boundary conditions between homogeneous layers are used to get unique solution of differential wave equation. These boundary equations form a system which can be solved for wave number as unknown. With the increasing number of layers in the waveguide the order of that system increases thus increasing its complexity and its solution becomes more complicated. Therefore, the time of computation rises dramatically either the accuracy of the solution decreases with the raise of the complexity of the structure. However, some papers describe waveguides with continuously inhomogeneous dielectric filling (Kiang, 1996) but only for axisymmetric modes.

MATERIALS AND METHODS

The proposed method uses Green's functions method for dispersion equations composing in cylindrical

radially piecewise homogeneous waveguides. This method is described by Knyazev *et al.* (2011) and is successfully implemented by Abdullin and Mitelman (2014) and (Abdullin *et al.*, 2014) for radiation and propagation problems in rectangular slotted waveguide.

The idea is based on the Fourier integral-series spatial decomposition of the electric and magnetic field's components in the cylindrical structures:

$$E_z = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{-jm\varphi} \int_{-\infty}^{\infty} E_{zmh}(r) e^{-jh z} dh \quad (1)$$

$$H_z = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{-jm\varphi} \int_{-\infty}^{\infty} H_{zmh}(r) e^{-jh z} dh \quad (2)$$

Where:

m = The discrete azimuthal index of mode

h = The continuous wavenumber

Describing electromagnetic field as the sum of electrical and magnetic waves propagating along z direction we can convert wave equations to the system of two equations. For example for TM (E) waves:

$$\begin{aligned} -\frac{d}{dr}(E_{zmh}^e) &= -j \frac{\gamma^2 Z_0}{\epsilon r k_0^2} (k_0 r H_{\varphi mh}^e) - \frac{h Z_0}{\epsilon k_0} J_{rmh}^{E\ ex} - J_{\varphi mh}^{M\ ex} \\ \frac{d}{dr}(k_0 r H_{\varphi mh}^e) &= j \frac{\epsilon k_0^2}{Z_0 \gamma^2} r \left[k^2 - h^2 - \left(\frac{m}{r} \right)^2 \right] E_{zmh}^e + \\ &k_0 r J_{zmh}^{E\ ex} - \frac{m h k_0}{\gamma^2} J_{\varphi mh}^{E\ ex} + \frac{\epsilon k_0}{\gamma^2 Z_0} m J_{rmh}^{M\ ex} \end{aligned} \quad (3)$$

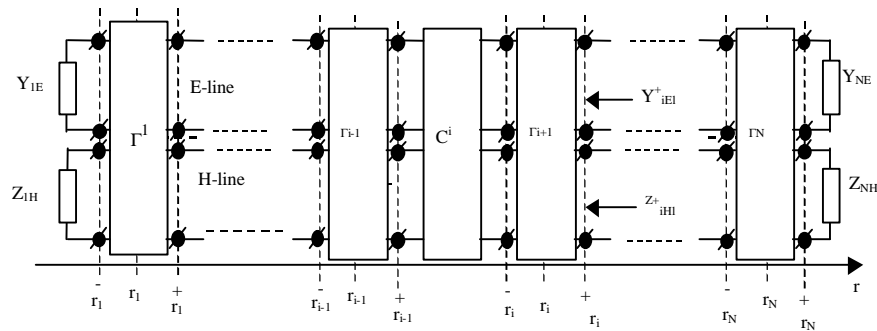


Fig. 1: Equivalent circuit of the radial lines

Where, $H_{qmh}^e, E_{zmh}^e, J_{zmh}^{E\text{ ex}}, J_{qmh}^{E\text{ ex}}, J_{zmh}^{M\text{ ex}}, J_{qmh}^{M\text{ ex}}$ are spectral components of the electric and magnetic fields, electric and magnetic external currents respectively, Y - transverse wave number. With interchange, $v_E = E_{zmh}^e, I_E = -k_0 r H_{qmh}^e$, we can regress Eq. 3 to system of telegraph equations:

$$\begin{aligned} -\frac{d}{dr} V_E &= jZ_E \chi I_E + v_{\text{ex}}^E \\ -\frac{d}{dr} I_E &= jY_E \chi V_E + i_{\text{ex}}^E \end{aligned} \quad (4)$$

Where, V_E, I_E, Z_E, Y_E are voltage, current, impedance and admittance of the equivalent radial E-line, v_{ex} and i_{ex} are equivalent voltage and current of sources $\chi = \sqrt{k^2 + h^2 - (m/r)^2}$ is the propagation constant in the equivalent radial transmission line. The same equations we can obtain for magnetic H-lines. Thus using the methods of circuits theory we can replace inhomogeneous structure with equivalent circuit consisting of two equivalent radial lines (Fig. 1).

Boundaries and layers in this circuit can be characterized by their transmission matrices Γ and C respectively which connect the voltages and currents on their ports. In the same way, we can describe the ends of these lines using impedances Z_{IH}, Z_{NH} and admittances Y_{IE}, Y_{NE} corresponding to the internal and external layers of the structure.

The described model can be used for example for field calculations in every layer for eigen modes of the multilayered circular waveguide. In this study, we apply this model to a circular metal layered waveguide and calculate its dispersion characteristics.

RESULTS AND DISCUSSION

To obtain the propagation constant of the waveguide we can use cross resonance conditions at the section p of radial line formulated as equations with directional modal impedances and admittances (Walter, 1990):

$$\begin{aligned} Z_{pHl}(m, h) + Z_{pHr}(m, h) &= 0 \\ Y_{pEl}(m, h) + Y_{pEr}(m, h) &= 0 \end{aligned} \quad (5)$$

The first equation solution is suitable for TM (E) modes and the second is for TE (H) modes. One should solve these equations for h as unknown on a particular frequency.

Described method allows creating an algorithm for dispersion analysis of cylindrical transmission lines with theoretically unlimited number of layers. Created program is able to calculate cut-off frequencies of eigen waves and plots of the propagation constant versus frequency in the cylindrical metal waveguide with 12 layers. Adding a new layer will lead to a non-significant increasing of computational time because in that case one should add only two more two-port devices in the equivalent circuit. To evaluate the operating frequency range of the waveguide, usually the concept of the Fractional Band Width (FBW) coefficient (the ratio between the critical frequencies of the first higher mode and the fundamental mode) is used:

$$T = \frac{f_c^h}{f_c^e} \quad (6)$$

The following types of cylindrical waveguides: empty waveguide, uniformly filled waveguide, waveguide with dielectric rod with dielectric cylinder with an air gap in the filling, with parabolic dielectric profile were analyzed.

TE₁₁ (H_{11}) wave is the fundamental wave in a cylindrical uniformly filled waveguides (Fig. 2 and 3). FBW in empty and filled by dielectric waveguides is $t = 1.306$. When using uniform dielectric filling, wavelengths of all modes decrease by $\sqrt{\epsilon}$. It means that usage of uniformly filled dielectric is helpful for miniaturization of waveguides. In addition the phase velocity and impedance decrease. To avoid high attenuation due to the dielectric losses one needs to use an expensive high quality materials. The y-axis of all

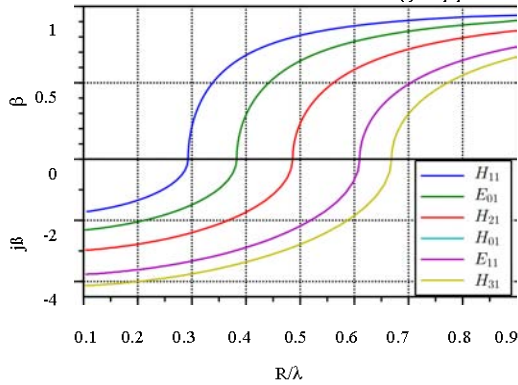


Fig. 2: Propagation constant of the eigen modes in the empty waveguide

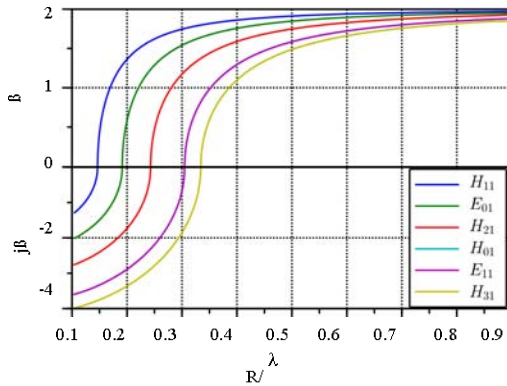


Fig. 3: Propagation constant of the eigen modes in the uniformly filled waveguide

frequency dependencies is normalized wave number. E_{01} mode is often being used in devices with rotating joint waveguides. Electromagnetic field must be uniform along the azimuthal coordinate. E_{01} wave is the first higher order mode in empty waveguide. To use this wave as operating one it is needed to suppress the fundamental wave with special mode filter. With using filters one can achieve FBW coefficient for this mode $t = 1.27$.

It is possible to create propagation conditions for E_{01} wave to make it the fundamental wave. To do this, a non-uniform dielectric filling can be used. Figure 4 and 5 show dependences for a waveguide with dielectric rod in the center. FBW coefficient becomes $t = 1.04$ (Fig. 2) without usage of special filters to suppress the HE_{11} wave. However, in the case we can suppress HE_{11} wave, coefficient increases to $t = 1.89$ and this value is quiet big for cylindrical waveguide. Using a larger value of permittivity (Fig. 5) increased FBW coefficient of the E_{01} mode as fundamental wave to $t = 1.4$ without filters and to $t = 1.97$ with suppressing of the hybrid mode HE_{11} .

The examples show that filling parameters could be chosen the way in which one can change the properties of the transmission line including fundamental mode.

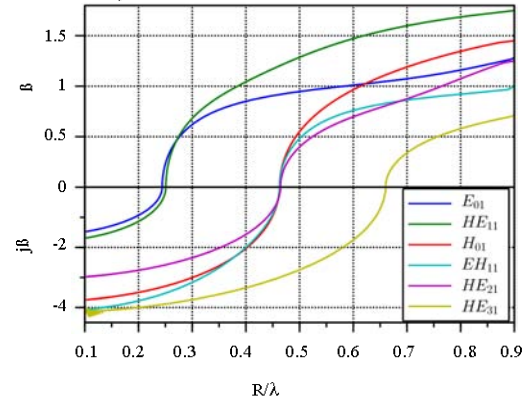


Fig. 4: Propagation constant of the eigen modes in the waveguide with dielectric rod ($\epsilon = 4$)

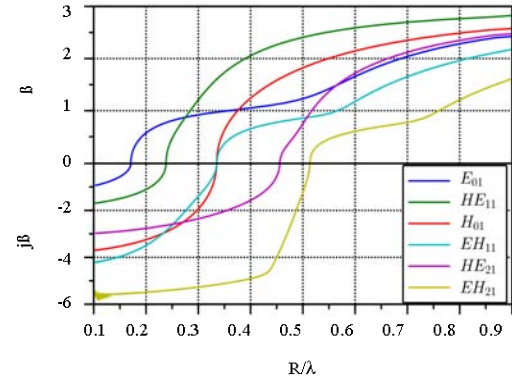


Fig. 5: Propagation constant of the eigen modes in the waveguide with dielectric rod ($\epsilon = 9$)

This effect was demonstrated in the past, i.e., in (Tsandoulas and Ince, 1971). In addition, dielectric rod allowed decreasing the dispersion for E_{01} and all EH_{m1} . Dispersion decreasing of fundamental wave mode allows transmitting signals with low distortions.

The use of H_{01} mode in circular waveguides is very attractive in terms of low attenuation. In the uniformly filled metal cylindrical waveguide H_{0n} and E_{n1} waves have the same cut-off frequencies and their frequency dependences of wave number are identical. This point makes it more difficult to use H_{01} mode because it is scattered by randomly located discontinuities in the waveguide and converted into a spectrum of waves with small amplitudes including E_{01} wave. Since, the speed of propagation of H_{01} and E_{11} modes are very close the amplitude of the E_{11} wave tends to accumulate when other amplitudes still small, so H_{01} wave constantly converts into E_{11} wave.

In inhomogeneous waveguides cut-off frequencies of H_{01} and EH_{11} modes (E_{11} becomes EH_{11}) are still equal but propagation constants become different. It creates conditions for more stable propagation of H_{01} mode.

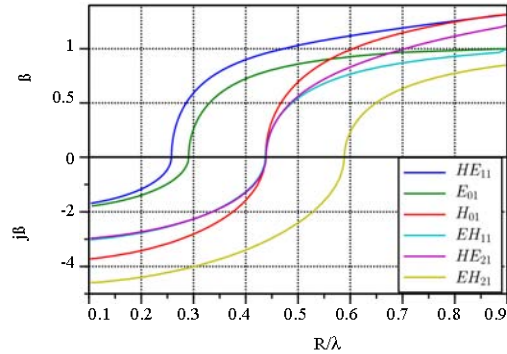


Fig. 6: Propagation constant of the modes in the waveguide with dielectric cylinder

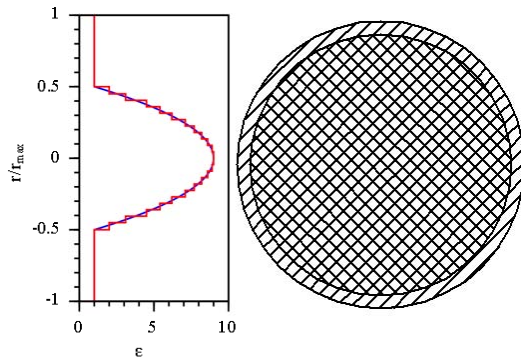


Fig. 7: Parabolic dielectric profile of the complex filled waveguide

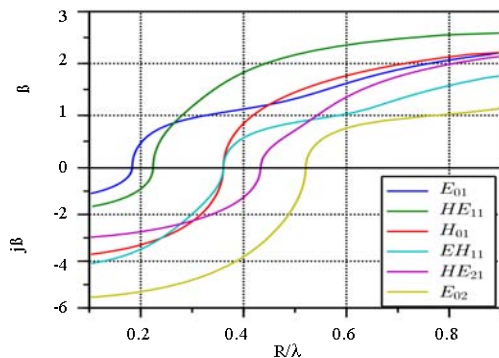


Fig. 8: Propagation constant of the modes in waveguide with parabolic dielectric profile

Splitting of H_{01} and EH_{11} is possible in all non-uniformly filled structures. The best result in splitting can be obtained using dielectric cylinder (Fig. 6).

Created software allows calculating characteristics of the structures with up to 12 layers. Non-uniform dielectric profile (Fig. 7) can be approximated by a stepped function, so the described method can be useful in stratified filling researches. The results of simulation are shown in Fig. 8.

CONCLUSION

In the study, the novel effective method for solving the problems of dispersion calculation of the circular metal waveguides with inhomogeneous filling is described. It is shown that calculation not only the phase constant but also an attenuation of cut-off modes and overmoded waveguides is possible. Continuous dielectric profiles can be approximated as accurately stratified media. In that case, the efficiency of the method is provided by its matrix nature and using of the recurrent equations of the circuit theory.

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